

Mixed-Port Scattering and Hybrid Parameters for High-Speed Differential Lines

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Abstract—High-speed transmission lines are commonly routed as differential lines to control sensitivity to noise on the reference planes at higher speeds. Differential lines are typically characterized in terms of mixed-mode scattering parameters, as they provide insight into the behavior of differential and common signals, as well as the mode conversion among them. These mixed-mode scattering parameters can be mathematically obtained from single-ended parameters, which can, for example, be measured with a four-port vector network analyzer. There has been recent efforts to develop extended or modified versions of mixed-mode scattering parameters, especially for tightly-coupled lines. This can be a point of confusion in interpreting the behavior of differential lines. In this paper, we introduce the mixed-port scattering and hybrid parameters, which do not suffer from any such ambiguous definitions. Mixed-port hybrid parameters are the most natural way to represent any four-port differential circuit, as they are based on intuitive differential and common-port excitations of the network. They also enable extraction of the current division factor experimentally, which is a critical parameter for electromagnetic interference analysis of differential lines. Mixed-port scattering parameters are also defined based on common and differential port excitations, allowing a simpler interpretation than their mixed-mode counterparts, without the need for defining even, odd, common, or differential-mode impedances. As such, mixed-port scattering and hybrid parameters can be used to analyze the performance of a general differential network, certainly including coupled or asymmetrical lines, without any ambiguity.

Index Terms—Current division factor, differential lines, mixed-mode scattering parameters, mode conversion.

I. INTRODUCTION

SINGLE-ENDED to mixed-mode conversion can be considered as a linear transformation of nodal voltages and currents into modal ones, by assigning a differential and common-mode voltage and current to each port pair. For the more

Manuscript received April 10, 2018; revised May 10, 2018 and June 11, 2018; accepted July 9, 2018. This work was supported in part by the National Science Foundation under Grant No. 1408637 and in part by Alexander von Humboldt Foundation. (*Corresponding author: Arif Ege Engin.*)

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Digital Object Identifier 10.1109/TEMC.2018.2855040

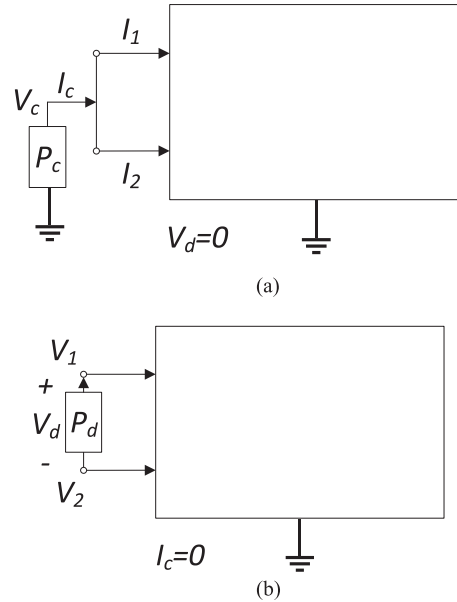


Fig. 1. New port definitions. (a) Common port (P_c). (b) Differential port (P_d).

commonly used mixed-mode scattering parameters, the transformation is from single-ended power waves into differential and common-mode power waves [1]. These linear transformations can be done with arbitrary coefficients, resulting in various definitions of mixed-mode parameters [2]–[4]. Mixed-mode parameters have also been defined for the impedance and admittance matrices [5], chain scattering parameters [6], as well as ABCD parameters [7]. A major advantage of working with mixed-mode parameters is that they allow to gain intuition into the properties of a differential network with a quick glance.

In this paper, we introduce *mixed-port* scattering and hybrid parameters, which provide an intuitive view into the definition of differential and common signals, while not suffering from ambiguities existing in earlier *mixed-mode* parameters. The new mixed-port parameters introduced in this paper are based on the common- and differential-port excitations as shown in Fig. 1. The mixed-port parameters obtained with such measurement ports will be referenced to the actual port impedances, without requiring the definition of even, odd, common, or differential-mode impedances.

Fig. 2(a) shows a pair of single-ended ports in a general microwave network. The traditional mixed-mode scattering parameters are calculated based on common-mode and

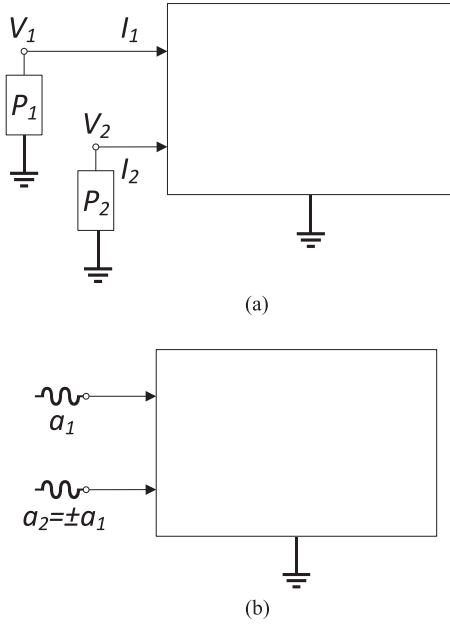


Fig. 2. Traditional port definitions for single-ended and mixed-mode excitation. (a) Single-ended pair of ports (P_1 and P_2). (b) Mixed-mode excitation by common ($a_2 = a_1$) or differential ($a_2 = -a_1$) mode power waves [1].

differential-mode power waves [1] as shown in Fig. 2(b). Such a measurement requires exciting both single-ended ports in-phase and with a 180° phase shift, which can, for example, be accomplished by using a balun. We eliminate the need for such phase alignments in the proposed mixed-port parameters.

To develop the intuitive nature of the mixed-port hybrid parameters, we start by defining possible port excitations for the differential and common signals. This allows us to relate the developed parameters to a possible simulation or measurement setup. The common port is across the shorted single-ended ports and ground reference as shown in Fig. 1(a). We define the differential port across the two single-ended ports as shown in Fig. 1(b). An advantage of such a mixed-port definition is that these ports can be realized in electromagnetic simulators (e.g., differential and common-mode ports in Sonnet [8]), and implemented on test boards similar to single-ended ports. Therefore, we define the ports in a more intuitive way as physical ports, as opposed to based on modes in traditional mixed-mode parameters. This eliminates the ambiguity in defining unequal even- and odd-mode impedances depending on the coupling level in differential lines [2], [9], [10]. In [11] and [12], we introduce the mixed-port hybrid and scattering parameters, respectively. In this paper, we present a comprehensive derivation of these parameters and their relation to each other. We also introduce a methodology to extract the current division factor experimentally; and provide simulation to hardware correlation. The advantages of these new mixed-port scattering and hybrid parameters over existing mixed-mode definitions can be summarized as follows.

- 1) The mixed-port hybrid parameters allow experimental characterization of the current division factor, which is a critical parameter to assess the behavior of differential lines for electromagnetic interference.

- 2) The mixed-port scattering parameters can be characterized with simple two-port measurement or simulation setups for the differential ports, common ports, or a combination of them. The reference impedance of such scattering parameters is unambiguously defined from the port impedances.

II. MIXED-PORT HYBRID PARAMETERS

A. Modal Voltages and Currents

The traditional definition of differential and common-mode voltages and currents are

$$\begin{pmatrix} V_d \\ V_c \end{pmatrix} = \begin{pmatrix} 1 & -1 \\ 1/2 & 1/2 \end{pmatrix} \begin{pmatrix} V_1 \\ V_2 \end{pmatrix} \quad (1)$$

and

$$\begin{pmatrix} I_d \\ I_c \end{pmatrix} = \begin{pmatrix} 1/2 & -1/2 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} I_1 \\ I_2 \end{pmatrix} \quad (2)$$

where the subscripts 1 and 2 represent the pair of single-ended ports. From these modal voltages and currents, differential and common-mode power waves can also be defined, which result in mixed-mode scattering parameters. In contrast, mixed-port hybrid parameters relate voltages and currents rather than power waves. Specifically, the common current is excited by defining a port across the shorted single-ended ports and ground reference, which results in $V_d = V_1 - V_2 = 0$, as shown in Fig. 1(a). We excite the differential voltage by defining a port across the two single-ended ports as in Fig. 1(b). This differential excitation results in $I_c = I_1 + I_2 = 0$. As a result, we obtain a straightforward way to enforce either I_c or V_d to zero, which results in mixed-port hybrid parameters.

B. Four-Port Networks

An important special case is when the network is a four-port, such as in a differential line. We define the mixed-port hybrid parameters in this case as

$$\begin{bmatrix} I_{d1} \\ I_{d2} \\ V_{c1} \\ V_{c2} \end{bmatrix} = \begin{bmatrix} G_{dd11} & G_{dd12} & G_{dc11} & G_{dc12} \\ G_{dd21} & G_{dd22} & G_{dc21} & G_{dc22} \\ G_{cd11} & G_{cd12} & G_{cc11} & G_{cc12} \\ G_{cd21} & G_{cd22} & G_{cc21} & G_{cc22} \end{bmatrix} \begin{bmatrix} V_{d1} \\ V_{d2} \\ I_{c1} \\ I_{c2} \end{bmatrix} \quad (3)$$

or in more compact form using submatrices as

$$\begin{bmatrix} I_d \\ V_c \end{bmatrix} = \begin{bmatrix} G_{dd} & G_{dc} \\ G_{cd} & G_{cc} \end{bmatrix} \begin{bmatrix} V_d \\ I_c \end{bmatrix}. \quad (4)$$

Since we can enforce I_c or V_d terms on the right-hand side to be equal to zero by differential and common-port excitations, the mixed-port hybrid parameters present a new intuitive approach for characterization of differential networks. The submatrices along the diagonal (G_{dd} and G_{cc} terms) can be calculated by two separate two-port measurements as shown in Fig. 3. The port configuration in Fig. 3(a) provides the elements of the G_{cc} matrix as $V_{d1} = V_{d2} = 0$ with the common-port excitation on

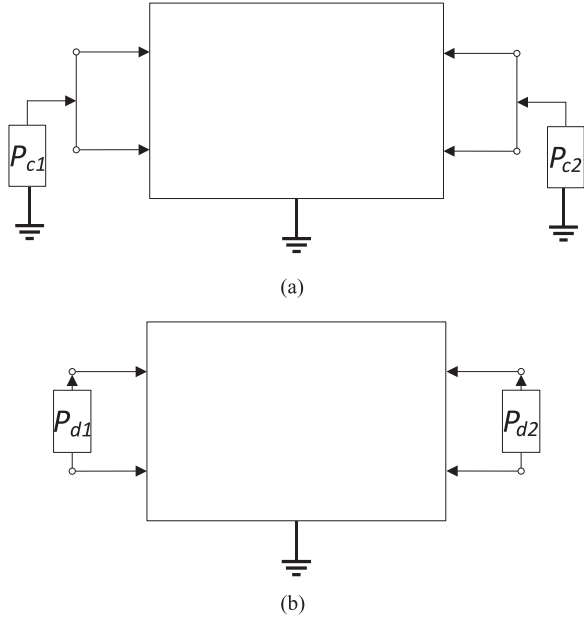


Fig. 3. Two-port excitation of a four-port network to extract. (a) Common. (b) Differential-port parameters.

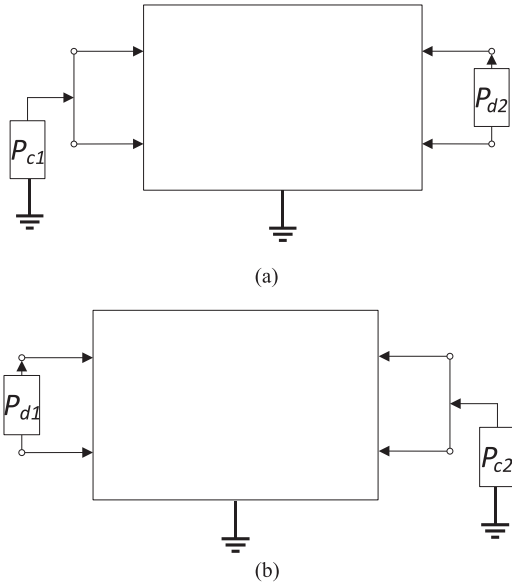


Fig. 4. Two-port excitation of a four-port network to characterize mode conversion.

both sides. From the port configuration in Fig. 3(b), we can extract the elements of G_{dd} .

Mode conversion can be characterized using the mixed-port excitations in Fig. 4. From Fig. 4(a), we can extract

$$\begin{bmatrix} G_{dd22} & G_{dc21} \\ G_{cd12} & G_{cc11} \end{bmatrix}. \quad (5)$$

Finally, from the port configuration in Fig. 4(b), we obtain

$$\begin{bmatrix} G_{dd11} & G_{dc12} \\ G_{cd21} & G_{cc22} \end{bmatrix}. \quad (6)$$

Some redundancy in these measurements can be observed. As a result, not every element of the hybrid mixed-mode parameters can be characterized with this measurement setup. Specifically, G_{dc11} , G_{dc22} , G_{cd11} , G_{cd22} parameters cannot be characterized with these four measurements.

Obtaining all the elements of mixed-port hybrid parameters would require additional measurements or simulations. One approach to obtain the full matrix would be conversion from other mixed-mode representations. For example, if the mixed-mode impedance matrix is given as

$$\begin{bmatrix} V_d \\ V_c \end{bmatrix} = \begin{bmatrix} Z_{dd} & Z_{dc} \\ Z_{cd} & Z_{cc} \end{bmatrix} \begin{bmatrix} I_d \\ I_c \end{bmatrix} \quad (7)$$

the hybrid parameters can be obtained as

$$\begin{bmatrix} I_d \\ V_c \end{bmatrix} = \begin{bmatrix} Z_{dd}^{-1} & -Z_{dd}^{-1}Z_{dc} \\ Z_{cd}Z_{dd}^{-1} & Z_{cc} - Z_{cd}Z_{dd}^{-1}Z_{dc} \end{bmatrix} \begin{bmatrix} V_d \\ I_c \end{bmatrix}. \quad (8)$$

Single-ended impedance parameters, given by $V = ZI$, can also be directly converted to mixed-port hybrid parameters as

$$G = \frac{1}{2}(M_1Z + M_2)(M_2Z + M_1)^{-1} \quad (9)$$

based on the transformation matrices M_1 and M_2 defined as

$$M_1 = \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 1 & 0 & 0 \\ 0 & 0 & 1 & 1 \end{bmatrix} \quad (10)$$

$$M_2 = \begin{bmatrix} 1 & -1 & 0 & 0 \\ 0 & 0 & 1 & -1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \quad (11)$$

which transform single-ended ports to mixed ports

$$\begin{bmatrix} I_d \\ V_c \end{bmatrix} = \frac{1}{2}M_1V + \frac{1}{2}M_2I \quad (12)$$

$$\begin{bmatrix} V_d \\ I_c \end{bmatrix} = M_2V + M_1I. \quad (13)$$

The mixed-port hybrid matrix can also be completely characterized without conversion from other parameters. This is based on additional two-port measurements to fill in the missing four terms as it will be described in Section II-C.

C. Characterization of Current Division Factor

An important special case to consider is asymmetrical transmission lines. Traditional mixed-mode scattering parameters will exhibit mode conversion for this case. An extended mixed-mode scattering matrix is defined in [2], which results in no mode-conversion by making use of a current division factor to perform the modal transformations. The new differential-port parameters in this paper will be the same as the extended parameters for such asymmetrical transmission lines. An advantage of the new differential-port parameters is that it does not require

a prior calculation of the current division factor. As such, it retains the simplicity of the traditional mixed-mode scattering parameters.

The current division factor h in [2], [13] is calculated from the transmission line inductance or capacitance matrix. In this paper, we introduce a measurement based characterization of h . This parameter can be extracted from G_{dc11} for the left side, or G_{dc22} for the right side. For the example of left side, G_{dc11} is defined as

$$G_{dc11} = \frac{I_{d1}}{I_{c1}} = \frac{1}{2} \frac{I_1 - I_2}{I_1 + I_2}. \quad (14)$$

The current division factor is defined as the ratio of the current in port 1 to the total common-mode current. Based on the common port excitation in Fig. 1(a), h can be calculated as

$$\begin{aligned} h &= \frac{I_1}{I_{c1}} = \frac{I_1}{I_1 + I_2} = \frac{1}{2} \left(1 + \frac{I_1 - I_2}{I_1 + I_2} \right) \\ &= \frac{1}{2} (1 + 2G_{dc11}) = \frac{1}{2} + G_{dc11}. \end{aligned} \quad (15)$$

Unlike the transmission-line approach in [2], this current division factor is not limited to differential lines and can be calculated for an arbitrary microwave network. The current division factor can be different for the left and right sides on the network. Equivalent voltage division factors can also be defined based on G_{cd11} and G_{cd22} parameters. For the example of the left side of the network, we define the voltage division factor as the ratio of the negative of the voltage in port 2 to the differential-port voltage, which can be obtained as

$$h = \frac{-V_2}{V_1 - V_2} = \frac{1}{2} \left(1 - \frac{V_1 + V_2}{V_1 - V_2} \right) = \frac{1}{2} (1 - 2G_{cd11}) = \frac{1}{2} - G_{cd11}. \quad (16)$$

Both h factors in (15) and (16) are equivalent (i.e., $G_{dc11} = -G_{cd11}$) for reciprocal networks. To prove that, consider the single-ended excitation of the two-ports on the left side of the network as shown in Fig. 2(a). Enforcing the differential voltage and common current to zero on the right side of the network would require short circuiting the two right ports together, which is not shown in the figure. Assume now that the impedance matrix for the two ports on the left side of the network is obtained as

$$\begin{bmatrix} V_1 \\ V_2 \end{bmatrix} = \begin{bmatrix} Z_{11} & Z_{12} \\ Z_{21} & Z_{22} \end{bmatrix} \begin{bmatrix} I_1 \\ I_2 \end{bmatrix}. \quad (17)$$

This single-ended impedance matrix can be used to calculate the current and voltage division factors. The current division factor is defined when there is a common-port excitation, or $V_1 = V_2$. Enforcing this condition on the single-ended impedance matrix results in

$$Z_{11}I_1 + Z_{12}I_2 = Z_{21}I_1 + Z_{22}I_2 \quad (18)$$

which yields the current division factor in (15) as

$$h = \frac{I_1}{I_1 + I_2} = \frac{Z_{22} - Z_{12}}{Z_{11} - Z_{12} - Z_{21} + Z_{22}}. \quad (19)$$

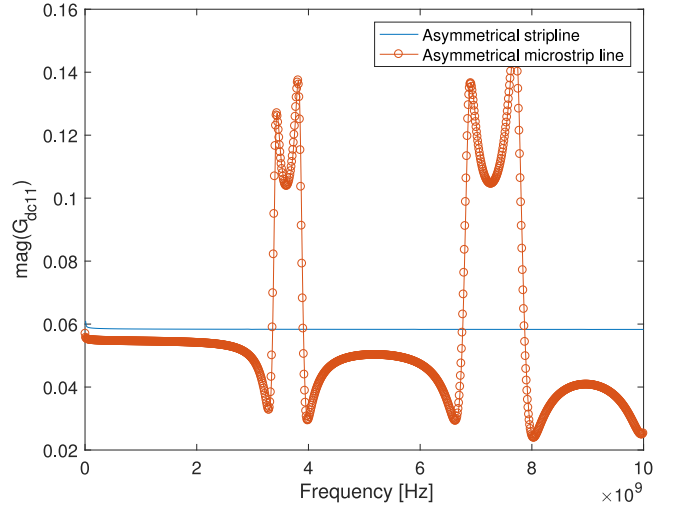


Fig. 5. $G_{dc11} = I_{d1}/I_{c1}$ represents the ratio of differential current to common current. Asymmetrical microstrip lines result in strong-frequency dependent variation due to the presence of inhomogeneous medium.

In a similar way, the voltage division factor in (16) is obtained when there is a differential-port excitation, or $I_1 = -I_2$. Enforcing this condition results in

$$V_1 = (Z_{11} - Z_{12})I_1 \quad (20)$$

$$V_2 = (Z_{21} - Z_{22})I_1 \quad (21)$$

which yields the voltage division factor in (16) as

$$h = \frac{-V_2}{V_1 - V_2} = \frac{Z_{22} - Z_{21}}{Z_{11} - Z_{12} - Z_{21} + Z_{22}}. \quad (22)$$

This shows that the current and voltage division factors are equivalent for any reciprocal network. Based on these additional single-ended two-port measurements, current or voltage division factors can therefore be calculated to obtain the remaining four elements G_{dc11} , G_{dc22} , G_{cd11} , G_{cd22} .

Current or voltage division factors can vary as a function of frequency for asymmetrical transmission lines in an inhomogeneous medium (e.g., microstrip lines), as the transformation matrices given in [2] do not diagonalize the inductance and capacitance matrices simultaneously for the case of inhomogeneous medium [14]. As an example, a 1-in long coupled microstrip line is simulated using Keysight ADS [15], on a substrate with a thickness of $100 \mu\text{m}$, dielectric constant of 4, and loss tangent of 0.025. The conductors are 1 oz. copper, with a spacing of $100 \mu\text{m}$ between line 1 and line 2, which have a width of $150 \mu\text{m}$ and $100 \mu\text{m}$, respectively. Another simulation was done by covering the microstrip lines with another $100 \mu\text{m}$ -thick substrate and a ground plane to obtain a coupled stripline configuration. Fig. 5 shows the variation of G_{dc11} as a function of frequency. A constant current division factor h , or G_{dc11} , is obtained for the stripline case. The microstrip line, on the other hand, has strong-frequency dependent variation due to the presence of inhomogeneous medium. As a result, the extended mixed-mode parameters in [2] would not decouple the differential and common modes in this case.

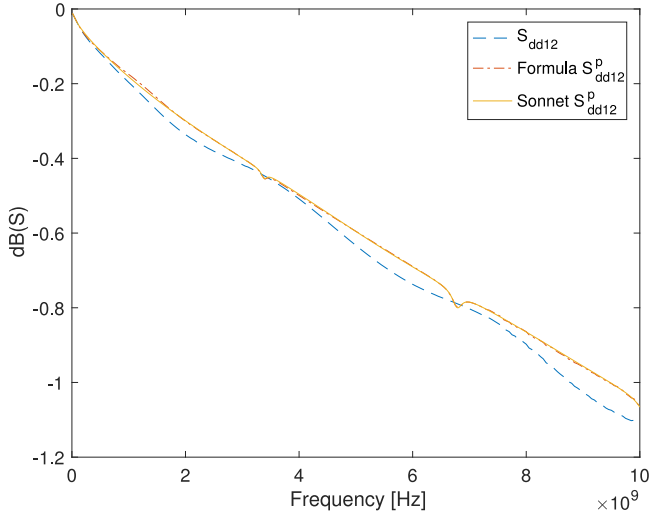


Fig. 6. Traditional differential-mode scattering parameters (S_{dd}) can be different than new differential-port scattering parameters (S_{dd}^p) for asymmetrical lines. The Sonnet simulation setup corresponds to the differential-port setup in Fig. 3(b); therefore, it matches the calculated S_{dd}^p from the formula in (24).

III. MIXED-PORT SCATTERING PARAMETERS

For high speed or microwave networks, the measurement setups in Figs. 3 and 4 can be used to characterize the scattering parameters. These scattering parameters are different than the traditional mixed-mode scattering parameters. For example, the differential-port scattering parameters in Fig. 3(b) are obtained under the condition that the common port is left unterminated. Some differential signal interface standards such as low-voltage differential signaling (LVDS) indeed specify a single-resistor (differential) termination [16], [17]. In addition, eye diagrams of differential lines are occasionally generated with a solely differential termination [18], [19]. The new mixed-port scattering parameters would be a more intuitive definition of scattering parameters in such cases where only the differential mode is terminated. In contrast, traditional differential-mode scattering parameters assume that there is a termination of common-mode signals as well.

The traditional mixed-mode scattering parameters are given as

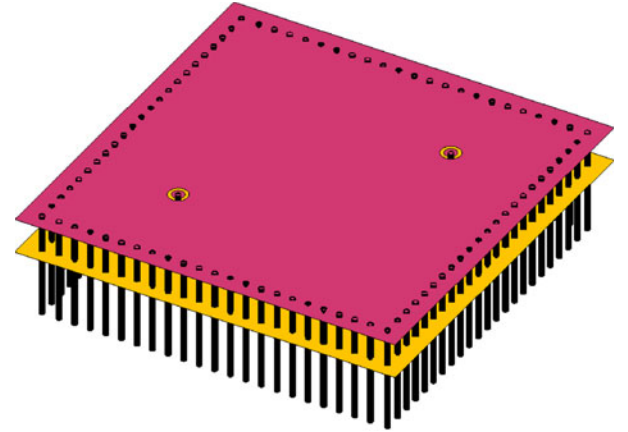
$$\begin{bmatrix} b_d \\ b_c \end{bmatrix} = \begin{bmatrix} S_{dd} & S_{dc} \\ S_{cd} & S_{cc} \end{bmatrix} \begin{bmatrix} a_d \\ a_c \end{bmatrix} \quad (23)$$

where S_{dd} is a 2×2 matrix representing the differential-mode scattering parameters. The conversion from the traditional mixed-mode (S_{dd}) to new differential-port scattering parameters (S_{dd}^p) can then be obtained by assigning $I_c = 0$ or $b_c = a_c$ as

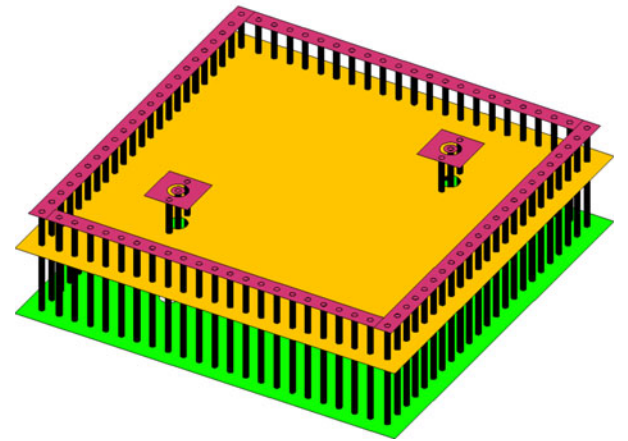
$$S_{dd}^p = S_{dd} + S_{dc} (1 - S_{cc})^{-1} S_{cd} \quad (24)$$

where S_{dd}^p is a 2×2 matrix with elements given as

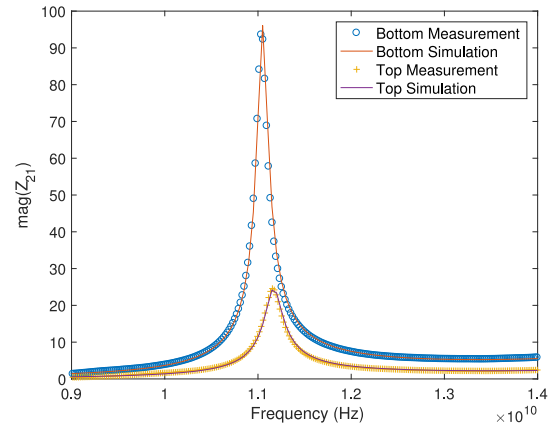
$$S_{dd}^p = \begin{bmatrix} S_{dd11}^p & S_{dd12}^p \\ S_{dd21}^p & S_{dd22}^p \end{bmatrix}. \quad (25)$$



(a)



(b)



(c)

Fig. 7. Shorted cavity resonators used to extract the dielectric thickness, loss tangent, and dielectric thickness. (a) Top substrate. (b) Bottom substrate. (c) Excellent match between simulation and measurement using the extracted parameters.

As expected, the two scattering parameters will be the same for cases where there is no mode conversion (i.e., $S_{dc} = 0$ or $S_{cd} = 0$ as in a symmetrical differential line.) S_{dd}^p is obtained from the two-port measurement setup in Fig. 3(b), and has a reference impedance equal to the port impedances in the usual way. As such, there is no need to define modes or mode impedances,

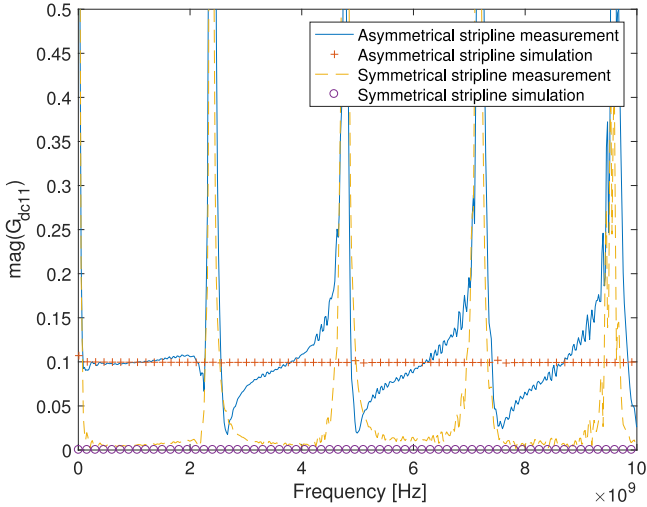


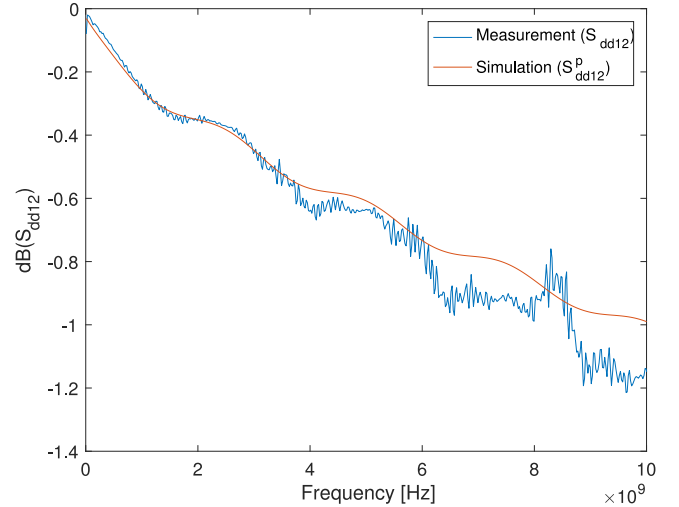
Fig. 8. $G_{dc11} = I_{d1}/I_{c1}$ represents the ratio of differential current to common current. Symmetrical lines will exhibit no mode conversion and have ideally $G_{dc11} = 0$. Simulation agrees well with simulation until the line length approaches 180° at around 2.5 GHz.

resulting in a more intuitive representation of differential circuits compared to traditional mixed-mode scattering parameters. The common-port scattering parameters S_{cc}^p will be extracted from the two-port measurement setup in Fig. 3(a). The mixed-port scattering parameters S_{cd}^p where the left port is the common port is obtained from the measurement in Fig. 4(a). Finally, the measurement setup in Fig. 4(b) would provide S_{dc}^p , where the left port is a differential port.

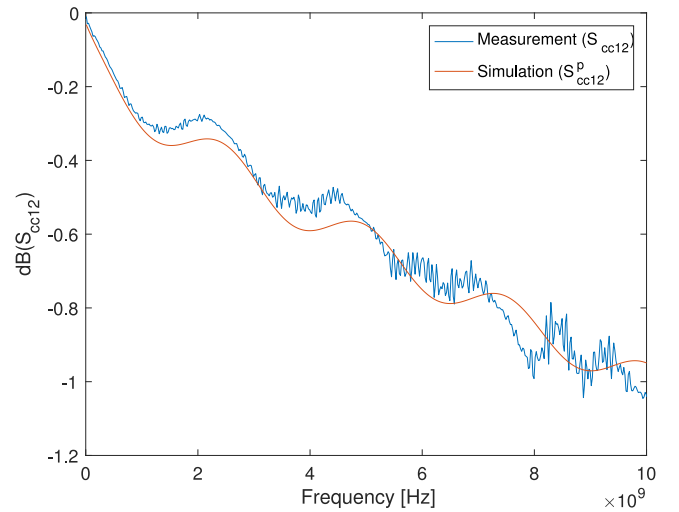
To provide a comparison between the traditional differential-mode scattering parameters (S_{dd}) and the new differential-port scattering parameters (S_{dd}^p), the same asymmetrical microstrip line example in Fig. 5 has been simulated with Sonnet. Traditional mixed-mode scattering parameters including S_{dd} are calculated from the four-port single-ended scattering parameters. Next, the same microstrip lines are simulated with differential ports in Sonnet. This simulation setup corresponds to the differential-port setup in Fig. 3(b); therefore, it matches the calculated S_{dd}^p from the formula in (24), as shown in Fig. 6. There is a very small discrepancy between calculated S_{dd}^p from formula in (24) and directly simulated S_{dd}^p from Sonnet, which may be due to any differences in deembedding of ports in the simulator.

IV. SIMULATION TO HARDWARE CORRELATION

Symmetrical and asymmetrical differential striplines have been characterized using a four-port vector network analyzer. The measured boards consisted of a prepreg as the top substrate of the stripline and a thicker bottom substrate made of MEGTRON6. The nominal thicknesses are $100 \mu\text{m}$ for the top substrate and $200 \mu\text{m}$ for the bottom one. The outer conductor layers are 1 oz. copper, whereas the inner layers are 0.5 oz. The nominal dielectric constant is 3.6 and loss tangent is 0.004 at 12 GHz according to the data sheet. To accommodate any deviations from the nominal values, it is important to extract the thickness, dielectric constant, as well as the loss tangent of the dielectrics. We designed simple shorted cavity resonators of



(a)



(b)

Fig. 9. Measurement to simulation correlation for symmetrical stripline. Mixed-port simulations are obtained using the port configurations in Fig. 3. (a) Differential insertion loss. (b) Common insertion loss.

size $10 \text{ mm} \times 10 \text{ mm}$ following the methodology in [20] and [21] as shown in Fig. 7.

The extracted parameters for the top substrate were thickness of $89 \mu\text{m}$, dielectric constant of 3.58, and loss tangent of 0.0094. For the bottom substrate, the same parameters were extracted as $211 \mu\text{m}$, 3.67, and 0.0069. The higher loss tangent compared to the nominal value can be attributed to the presence of any surface roughness that has not been accounted for during the extraction process. Using the extracted parameters, the shorted cavity resonators were simulated with the full-wave simulator Sonnet [8]. Fig. 7 shows excellent match between the simulation and measurement.

We apply the extracted parameters on the simulation of the striplines using Keysight ADS. The striplines had a length of 3.15 cm. In the symmetrical case, the width and spacing was $100 \mu\text{m}$; whereas the width of line 1 was increased to $200 \mu\text{m}$ for the asymmetrical case. Fig. 8 shows the variation of G_{dc11} as

a function of frequency. An approximately constant differential current to common current ratio is observed as expected. The deviation in measurements from this constant value at around 2.5 GHz can be attributed to the electrical length of the line being close to 180° . This common inaccuracy is observed in, for example TRL calibration, since the scattering parameters of the line standard becomes almost identical to a thru standard around those frequency points [22], and measurement inaccuracies due to the probing pads are amplified.

For symmetrical lines, there is no mode conversion, and the mixed-port scattering parameters are equivalent to the mixed-mode scattering parameters. The measured mixed-mode scattering parameters are compared against simulated mixed-port scattering parameters in Fig. 9. The good agreement indicates the accuracy of the extracted dielectric thickness, loss tangent, and dielectric thickness of the substrate layers. For asymmetrical lines, the new mixed-port parameters would be different than traditional mixed-mode parameters, due to the presence of mode conversion.

V. CONCLUSION

This paper introduced mixed-port scattering and hybrid parameters for differential circuits. These new parameters are based on simple port definitions, and do not make use of modal quantities such as even, odd, common, or differential-mode impedances. The new hybrid parameters are particularly useful to experimentally characterize the current division factor, a critical parameter for electromagnetic interference analysis of differential lines. The new mixed-port scattering parameters present a promising new way of characterizing and evaluating arbitrary differential circuits, without the ambiguities of existing modified or extended mixed-mode scattering parameter definitions.

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